

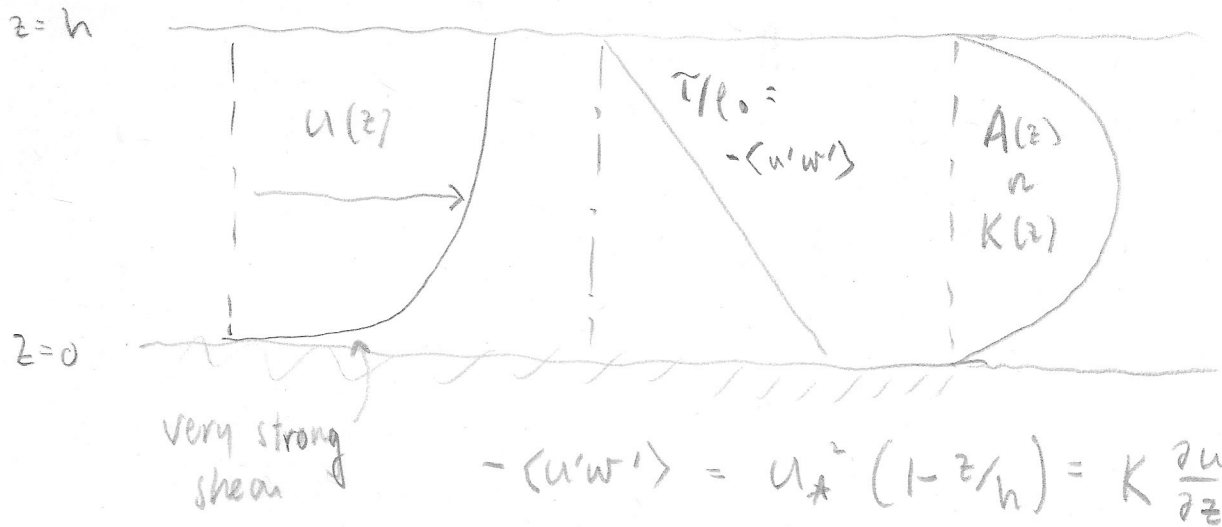
RG

"Phenomenology" of Boundary-Layer Turbulence

7/24/2014

(1)

- what we observe ; what we think matters

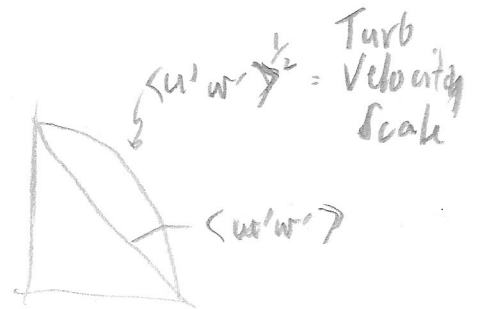


$$-\langle u'w' \rangle = u_*^2 (1 - z/h) = K \frac{\partial u}{\partial z}$$

You can't have a no slip condition without viscosity

$$K = L^2 T^{-1} = \text{velocity} \times \text{Length}$$

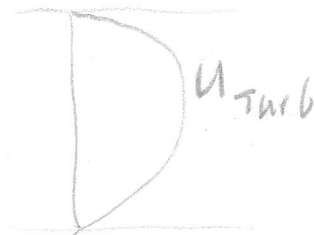
$$u_{\text{Turb}} \sim u_* \approx \langle u'w' \rangle^{1/2}$$

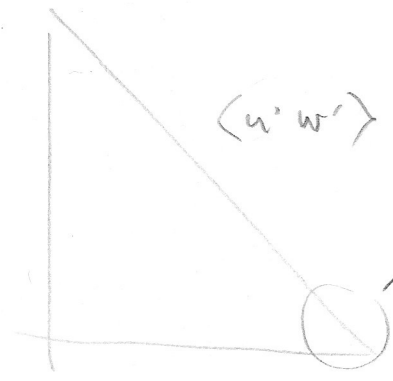


$$u_{\text{Turb}} \approx u_* \left(1 - \frac{z}{h}\right)^{1/2} \quad \text{and} \quad K = \kappa u_* \left(1 - \frac{z}{h}\right)$$

$$K(z) = u_{\text{Turb}} L_{\text{Turb}} \Rightarrow L_{\text{Turb}} = K / u_{\text{Turb}}$$

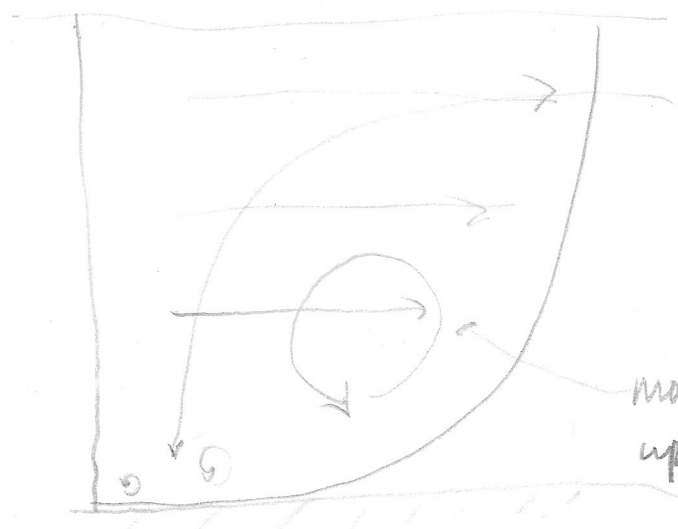
$$L_{\text{Turb}} = \kappa z \left(1 - \frac{z}{h}\right)^{1/2}$$





bottom 10% of z called
 "constant stress layer"
 and in it $\langle u'w' \rangle \approx u_*^2$

and log layer fit is most robust there
 (easy for other physics to affect things
 higher up.)

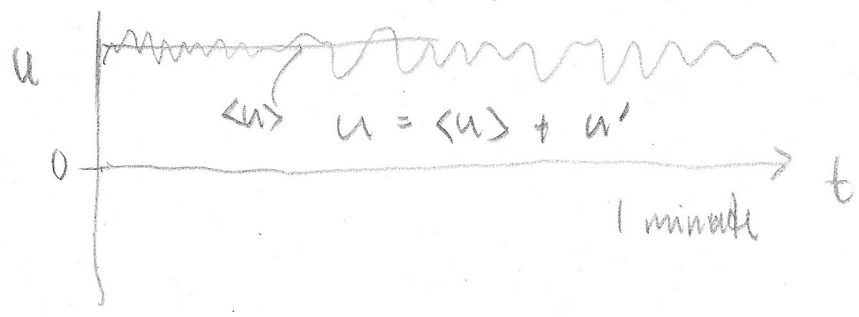


little eddies = small pressure
 perturbations force u that
 closes back on itself

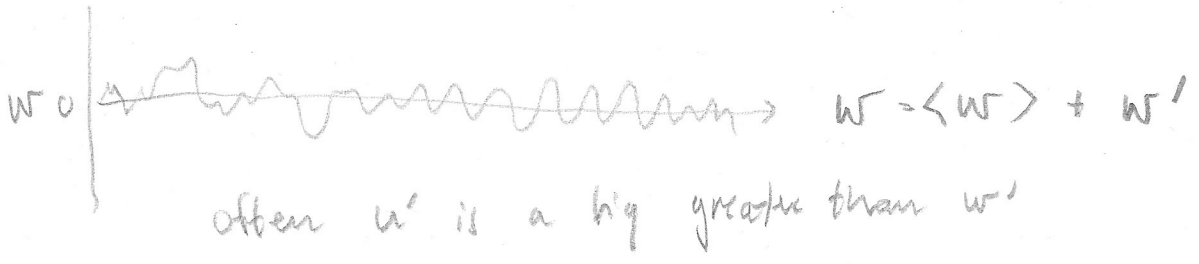
more room for bigger eddies
 up here

Turbulence ~ isotropic =

Turbulence Spectrum

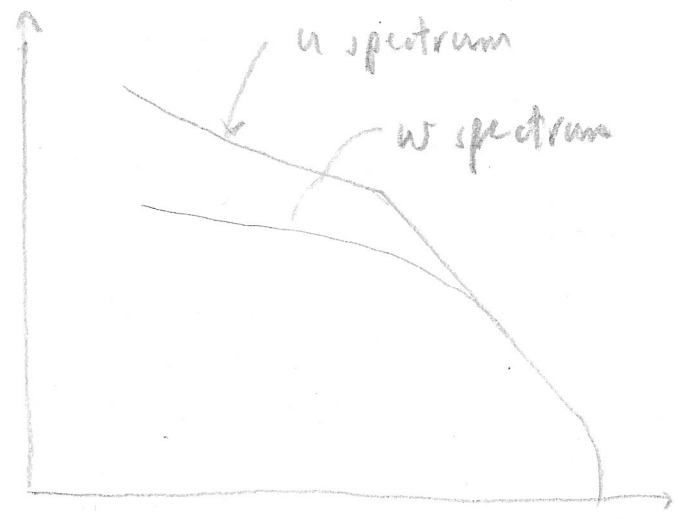


data from the Mast



often \$u'\$ is a bit greater than \$w'\$

$$u'(t) = \sum e^{i\omega t} \quad \text{FFT}$$



For your study place - what is the turbulent velocity scale, and where in the water column?

(extra credit - what is \$L_{turb}\$?) and what is \$\langle u \rangle\$?

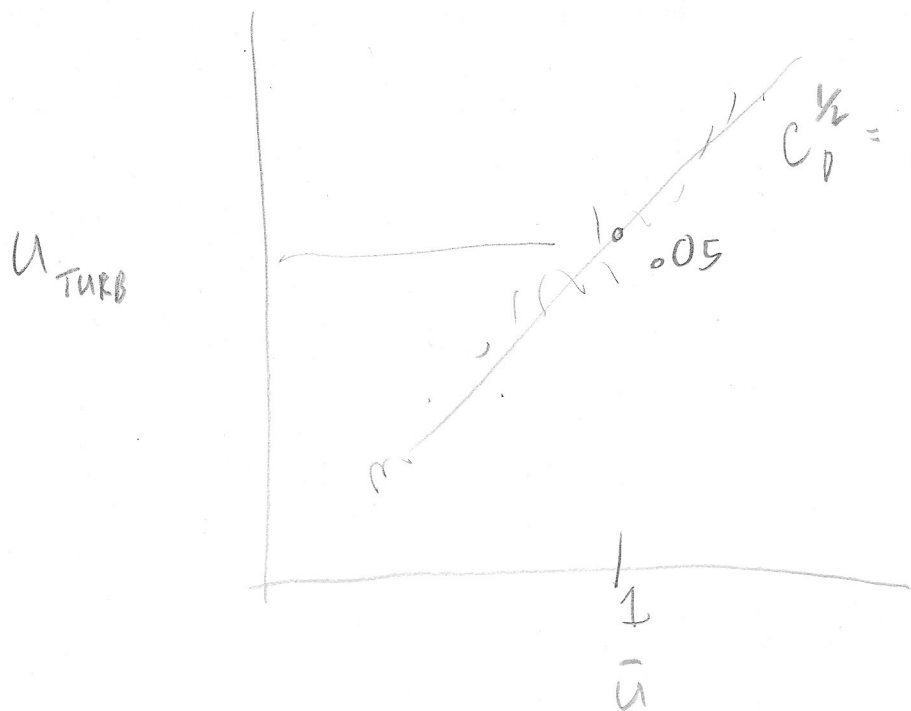
$$u'w' \sim K \frac{1u}{z^2} \sim C_d u^2 \sim 3 \times 10^{-3} \cdot 1^2 \quad \text{so } u_{turb} \leq \left(\frac{3 \times 10^{-3} \text{ m}^2}{s^2} \right)^{1/2} = 5 \text{ mm/sec!} = 5 \text{ cm/sec}$$

at 4 m \$L_{turb} \sim .4 \text{ m} \quad \langle u \rangle = 1 \text{ m/s}\$

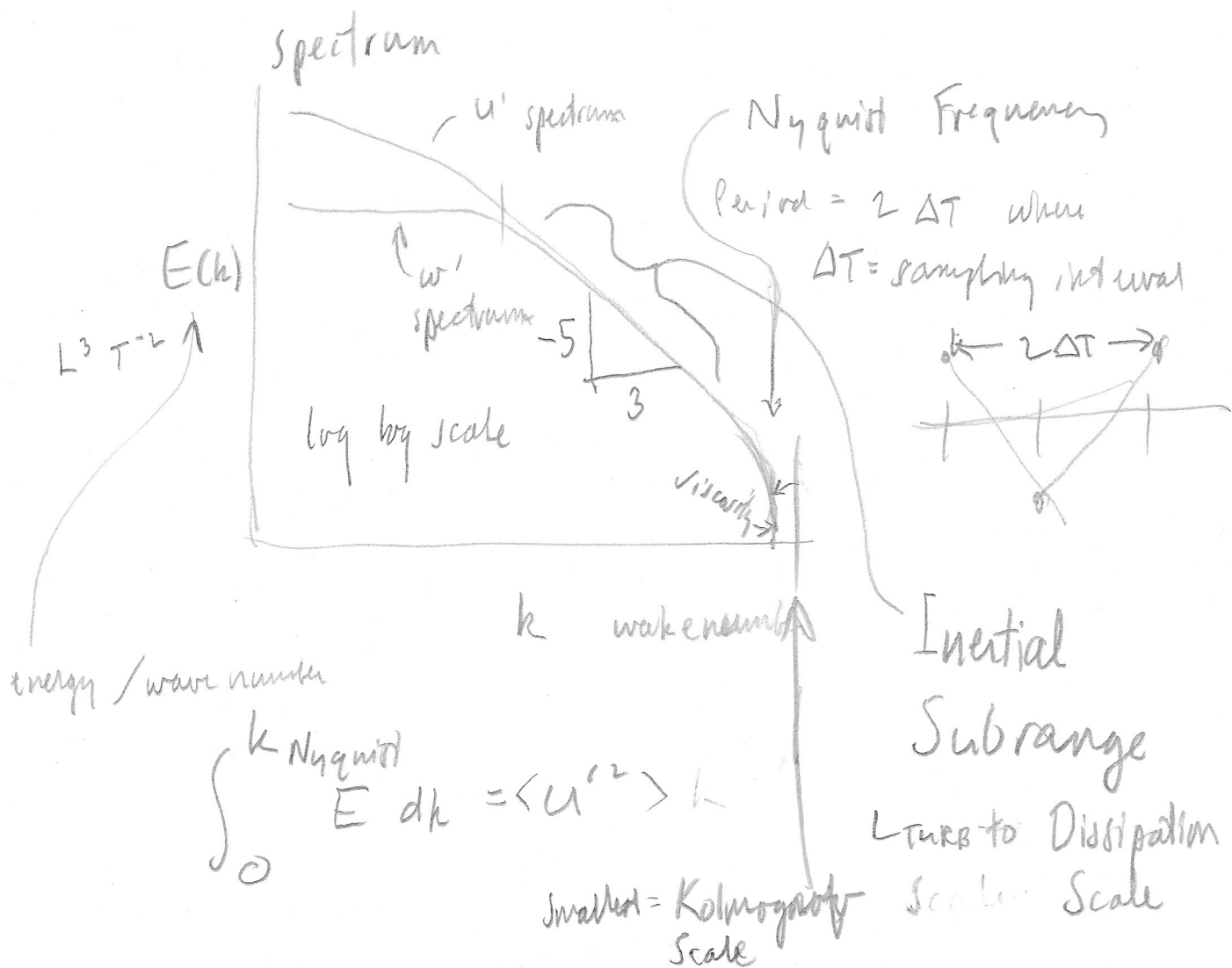
$\frac{3}{1000} \frac{1.5}{300} = \frac{1}{200}$

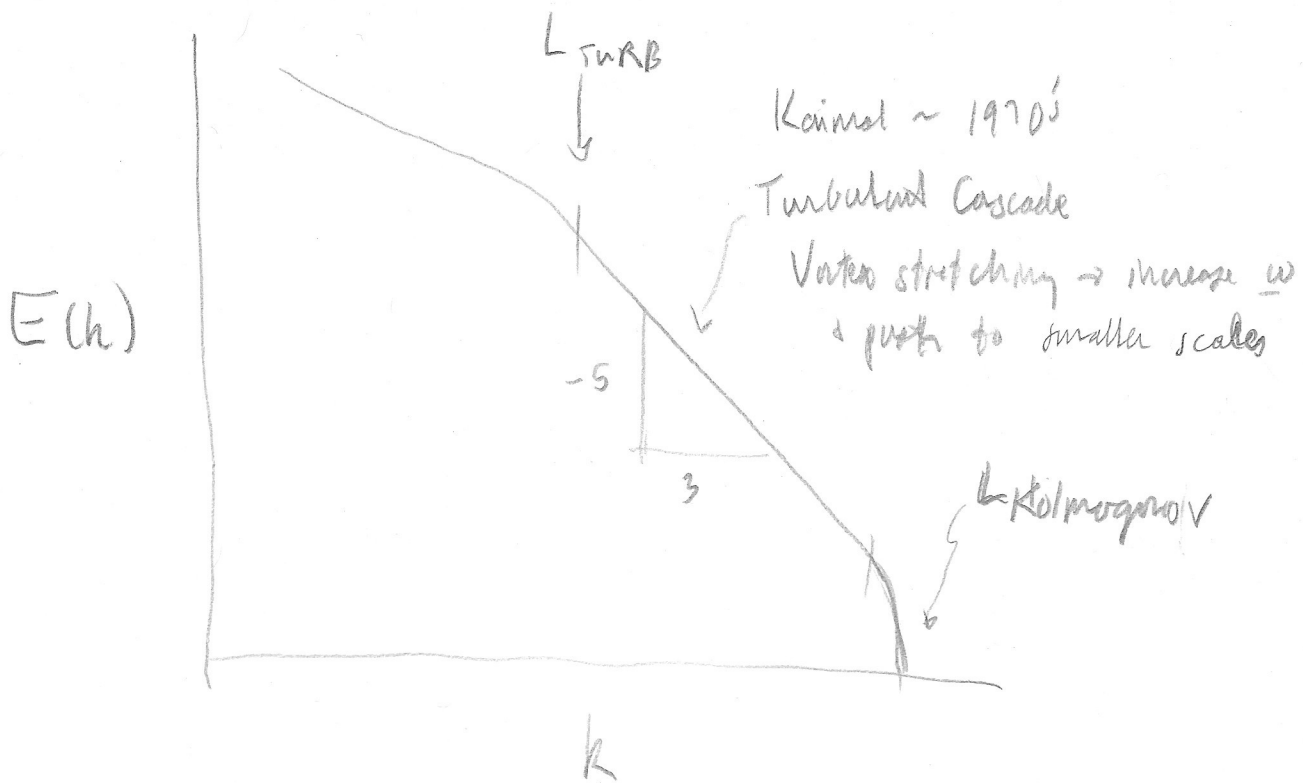
GT Tansla
Francis Brotherton
Peter Blinn
PM

(4)



$$u_{TURB} \sim C_D^{1/2} \bar{u}$$





- When eddies \rightarrow small scales it is a one way trip \rightarrow viscosity
- Generally there is creation of larger & smaller scales (where energy is lost)
 - \downarrow like billows in clouds

TKE equation

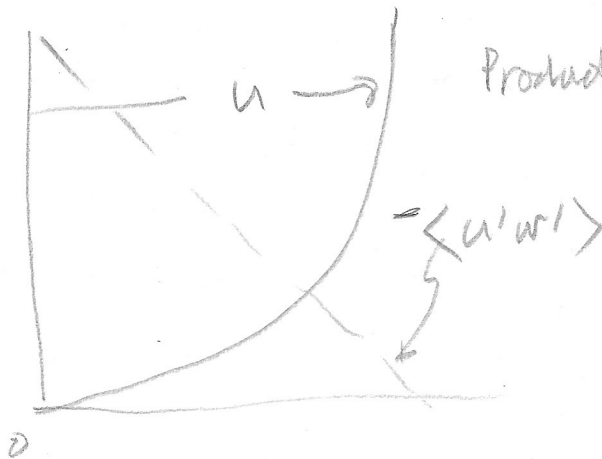
(6)

$$\left(\frac{\partial}{\partial t} + \underline{u} \cdot \nabla \right) TKE \equiv \text{Production} - \text{Dissipation}$$

Ignore compared

to RHS Terms \Rightarrow Production = Dissipation

What is Production?



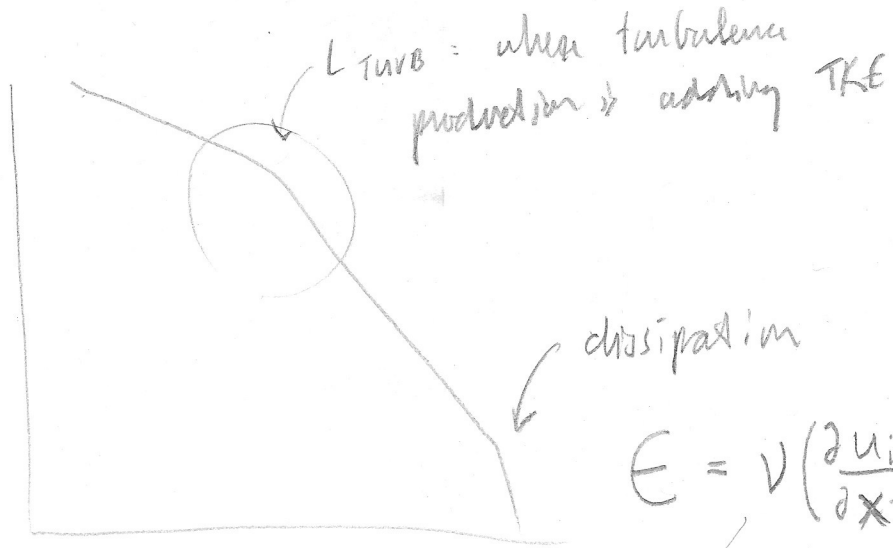
$$\text{Production} = -\langle u'w' \rangle \frac{\partial \langle u \rangle}{\partial z}$$

(Derived from an Energy Equation)

Loss of Mean KE
 \Rightarrow Gain of TKE

$$\left(\underline{u}_i \frac{\partial}{\partial x_i} \right) \langle u_s \rangle = \langle (u + u') \left[(u u)_x + (u v)_y + (u w)_z \right] \rangle$$

$$u \langle u'w' \rangle_z = \left[u \langle u'w' \rangle \right]_z - \langle u'w' \rangle \frac{\partial u}{\partial z}$$

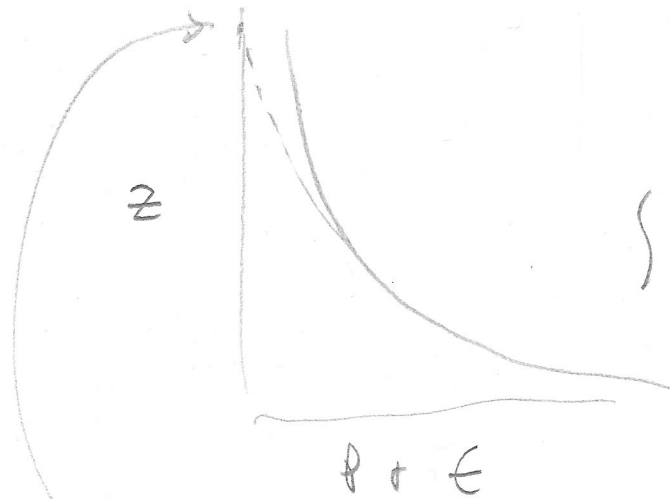


$$\epsilon = \nu \left(\frac{\partial u_i}{\partial x_j} \right)^2$$

$$\frac{L^2}{T} \frac{U}{T^2} \sim \frac{L^2}{T^3} \sim \frac{W}{kg}$$

in the log layer

$$\epsilon = - \underbrace{\langle u'w' \rangle}_{U_*^2} \underbrace{\frac{\partial \langle u \rangle}{\partial z}}_{\frac{U_*}{kz}} \quad (\text{Production})$$



al is finite because we start from z_0 instead of $z=0$. ✓

goes to zero at surface if you consider full expressions.

at 1 m a.b. $C_d = 3 \cdot 10^{-3}$, $z = 1 \text{ m}$, $u = 1 \text{ m/s}$

$$u_* = \sqrt{C_d} = 0.05$$

$$\epsilon = \frac{0.05^3}{0.4} = 3 \times 10^{-4} \frac{W}{kg}$$